

Normal completely positive maps on the space of quantum operations

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 - Dilation of quantum operations
- 2 Quantum supermaps
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 - Dilation of quantum supermaps
 - Examples
- 3 Superinstruments
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Mathematical setting

- $\mathcal{M} \subset \mathcal{L}(\mathcal{H}), \mathcal{N} \subset \mathcal{L}(\mathcal{K}) \dots$: complex and separable von Neumann algebras
- $\mathcal{M}_+, \mathcal{N}_+, \dots$: cones of their positive elements
- $M_n := \mathcal{L}(\mathbb{C}^n)$: von Neumann algebra of complex $n \times n$ -matrices
- $\mathcal{M} \hat{\otimes} \mathcal{N}$: algebraic tensor product of \mathcal{M} and \mathcal{N}
- $\mathcal{M} \bar{\otimes} \mathcal{N}$: von Neumann algebra tensor product of \mathcal{M} and \mathcal{N}

Recall that

$$\mathcal{M} \hat{\otimes} \mathcal{N} \subset \mathcal{L}(\mathcal{H} \otimes \mathcal{K})$$

and

$$\mathcal{M} \bar{\otimes} \mathcal{N} = \text{weak*}-\text{closure of } \mathcal{M} \hat{\otimes} \mathcal{N}$$

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Remarks

- $\mathcal{L}(\mathcal{H}) \bar{\otimes} \mathcal{L}(\mathcal{K}) = \mathcal{L}(\mathcal{H} \otimes \mathcal{K})$
- If $\dim \mathcal{N} < \infty$, then $\mathcal{M} \hat{\otimes} \mathcal{N} = \mathcal{M} \bar{\otimes} \mathcal{N}$
- If $\dim \mathcal{N}_i < \infty$ and $\mathcal{E} : \mathcal{M}_1 \rightarrow \mathcal{M}_2$, $\mathcal{F} : \mathcal{N}_1 \rightarrow \mathcal{N}_2$ are linear maps, then $\mathcal{E} \otimes \mathcal{F} : \mathcal{M}_1 \bar{\otimes} \mathcal{N}_1 \rightarrow \mathcal{M}_2 \bar{\otimes} \mathcal{N}_2$ is well-defined

Definition

A linear map $\mathcal{E} : \mathcal{M} \rightarrow \mathcal{N}$ is *completely positive* if the map

$$\mathcal{E} \otimes \mathcal{I}_n : \mathcal{M} \bar{\otimes} M_n \rightarrow \mathcal{N} \bar{\otimes} M_n$$

is positive for all $n \in \mathbb{N}$.

Here, $\mathcal{I}_n : M_n \rightarrow M_n$ is the identity map

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Definition

A linear map $\mathcal{E} : \mathcal{M} \rightarrow \mathcal{N}$ is called **positive** if $\mathcal{E}(\mathcal{A}) \subseteq \mathcal{A}$ for every $\mathcal{A} \subseteq \mathcal{M}$ such that $\mathcal{A} \subseteq \mathcal{M}_+$. It then makes sense to speak about positivity and boundedness of $\mathcal{E} \otimes \mathcal{F}$.

$$\mathcal{E} \otimes \mathcal{I}_n : \mathcal{M} \otimes M_n \rightarrow \mathcal{N} \otimes M_n$$

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Definition of quantum operations and channels

Definition

A linear map $\mathcal{E} : \mathcal{M} \rightarrow \mathcal{N}$ is a *quantum operation* if it is

(i) completely positive

(ii) normal:

$\mathcal{E}(A_n) \uparrow \mathcal{E}(A)$ for all sequences $\{A_n\}_{n \in \mathbb{N}}$ in \mathcal{M}_+ s. t. $A_n \uparrow A$

(iii) subnormalized:

$$\mathcal{E}(I_{\mathcal{M}}) \leq I_{\mathcal{N}}$$

\mathcal{E} is a *quantum channel* if condition (iii) is replaced by

(iii') normalized:

$$\mathcal{E}(I_{\mathcal{M}}) = I_{\mathcal{N}}$$

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Notations

- $\text{CP}(\mathcal{M}, \mathcal{N})$: normal completely positive maps in $\text{Hom}_{\mathbb{C}}(\mathcal{M}, \mathcal{N})$
- $\text{CP}_0(\mathcal{M}, \mathcal{N})$: subset of quantum operations
- $\text{CP}_1(\mathcal{M}, \mathcal{N})$: subset of quantum channels

Clearly,

$$\text{CP}_1(\mathcal{M}, \mathcal{N}) \subset \text{CP}_0(\mathcal{M}, \mathcal{N}) \subset \text{CP}(\mathcal{M}, \mathcal{N})$$

(iii') normalized:

$$\mathcal{E}(I_{\mathcal{M}}) = I_{\mathcal{N}}$$

Stinespring Theorem

Theorem

Suppose $\mathcal{M} \subset \mathcal{L}(\mathcal{H})$, $\mathcal{N} \subset \mathcal{L}(\mathcal{K})$. A linear map $\mathcal{E} : \mathcal{M} \rightarrow \mathcal{N}$ is a quantum operation iff there exist a Hilbert space \mathcal{V} and a bounded operator $V : \mathcal{K} \rightarrow \mathcal{H} \otimes \mathcal{V}$, with $\|V\|_\infty \leq 1$, such that

$$\mathcal{E}(A) = V^*(A \otimes I_{\mathcal{V}})V \quad \forall A \in \mathcal{M}.$$

*In this case, \mathcal{E} is a quantum channel iff $V^*V = I_{\mathcal{K}}$*

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Statement of the problem

Problem

Characterize the transformations

$$S : \text{CP}_0(\mathcal{M}, \mathcal{N}) \rightarrow \text{CP}_0(\mathcal{M}, \mathcal{N})$$

or, more generally,

$$S : \text{CP}_0(\mathcal{M}_1, \mathcal{N}_1) \rightarrow \text{CP}_0(\mathcal{M}_2, \mathcal{N}_2)$$

which are admissible in Quantum Mechanics (*quantum supermaps*)

Applications

- Quantum information
- Quantum measurement theory
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Properties of supermaps: Linearity

A quantum supermap must preserve mixtures and $S(0) = 0$



S is convex and $S(0) = 0$



S uniquely extends to a linear map defined on $\text{span } \text{CP}_0(\mathcal{M}, \mathcal{N})$

But what is the linear space spanned by $\text{CP}_0(\mathcal{M}, \mathcal{N})$?

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Example

Suppose

$$\mathcal{M} = \mathcal{N} = \mathcal{L}(\mathcal{H}).$$

- If $\mathcal{H} = \mathbb{C}^n$

$$\begin{aligned}\text{span CP}_0(\mathcal{M}, \mathcal{N}) &= \text{span CP}_0(M_n, M_n) \\ &= M_{n^2} \\ &= \text{Hom}_{\mathbb{C}}(M_n, M_n)\end{aligned}$$

- If $\dim \mathcal{H} = \infty$

$$\text{span CP}_0(\mathcal{M}, \mathcal{N}) = ???$$

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$$\text{span CP}_0(\mathcal{M}, \mathcal{N}) = \text{span CP}_0(M_n, M_n)$$

By Choi isomorphism,

$$\text{CP}(M_n, M_n) \simeq M_{n^2} +$$

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Completely bounded maps: Definition

Definition

A linear map $\mathcal{E} : \mathcal{M} \rightarrow \mathcal{N}$ is *completely bounded* if $\exists C > 0$ such that

$$\|(\mathcal{E} \otimes \mathcal{I}_n)(A)\|_\infty \leq C \|A\|_\infty$$

for all $A \in \mathcal{M} \bar{\otimes} M_n$ and $n \in \mathbb{N}$

Example

Suppose $\mathcal{M} \subset \mathcal{L}(\mathcal{H})$, $\mathcal{N} = \mathcal{L}(\mathcal{K})$. If $E, F \in \mathcal{L}(\mathcal{K}; \mathcal{H})$, the map

$$\begin{aligned} E^* \odot_{\mathcal{M}} F : \mathcal{M} &\longrightarrow \mathcal{L}(\mathcal{K}) \\ A &\longmapsto E^* A F \end{aligned}$$

is in $\text{CB}(\mathcal{M}, \mathcal{L}(\mathcal{K}))$

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Notation

$\text{CB}(\mathcal{M}, \mathcal{N})$: completely bounded weak*-continuous maps from \mathcal{M} to \mathcal{N}

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Completely bounded maps: Properties

1 **Ordering:** $\text{CP}(\mathcal{M}, \mathcal{N})$ is a cone in $\text{CB}(\mathcal{M}, \mathcal{N})$, hence it induces a linear ordering in $\text{CB}(\mathcal{M}, \mathcal{N})$, which we denote by \preceq

2 **Tensoring:** If $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{N}_1)$, $\mathcal{F} \in \text{CB}(\mathcal{M}_2, \mathcal{N}_2)$, the product

$$\mathcal{E} \otimes \mathcal{F} : \mathcal{M}_1 \hat{\otimes} \mathcal{N}_1 \longrightarrow \mathcal{M}_2 \hat{\otimes} \mathcal{N}_2$$

uniquely extends to a map

$$\mathcal{E} \otimes \mathcal{F} \in \text{CB}(\mathcal{M}_1 \bar{\otimes} \mathcal{M}_2, \mathcal{N}_1 \bar{\otimes} \mathcal{N}_2)$$

3 **Spanning:** If $\mathcal{N} = \mathcal{L}(\mathcal{K})$, every $\mathcal{E} \in \text{CB}(\mathcal{M}, \mathcal{L}(\mathcal{K}))$ can be written

$$\mathcal{E} = \mathcal{E}_1 - \mathcal{E}_2 + i(\mathcal{E}_3 - \mathcal{E}_4).$$

for some $\mathcal{E}_1, \dots, \mathcal{E}_4 \in \text{CP}(\mathcal{M}, \mathcal{L}(\mathcal{K}))$. In particular,

$$\text{span } \text{CP}_0(\mathcal{M}, \mathcal{L}(\mathcal{K})) = \text{CB}(\mathcal{M}, \mathcal{L}(\mathcal{K}))$$

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Remark

If $\mathcal{M} \subset M_m$ and $\mathcal{N} \subset M_n$, then

$$\text{CB}(\mathcal{M}, \mathcal{N}) = \text{Hom}_{\mathbb{C}}(\mathcal{M}, \mathcal{N}).$$

Properties (1) and (2) are trivial, and (3) follows from Choi isomorphism

Remarks

- For composite systems,

$$\begin{aligned}\text{CB}(\mathcal{M} \bar{\otimes} M_m, \mathcal{N} \bar{\otimes} M_n) &= \text{CB}(\mathcal{M}, \mathcal{N}) \hat{\otimes} \text{Hom}_{\mathbb{C}}(M_m, M_n) \\ &= \text{CB}(\mathcal{M}, \mathcal{N}) \hat{\otimes} M_{mn}\end{aligned}$$

- If

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Properties of supermaps: Complete positivity

A quantum supermap must preserve quantum operations on composite systems

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Here, $I_n : \text{Hom}_{\mathbb{C}}(M_n, M_n) \rightarrow \text{Hom}_{\mathbb{C}}(M_n, M_n)$ is the identity map

In other words, S is completely positive if $S \otimes I_n$ preserves the linear ordering \preceq for all $n \in \mathbb{N}$

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A quantum supermap must preserve quantum operations on composite systems

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A linear map $S : \text{CB}(\mathcal{M}_1, \mathcal{N}_1) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{N}_2)$ is *completely positive* if

$$(S \otimes I_n)(\text{CP}(\mathcal{M}_1 \bar{\otimes} M_n, \mathcal{N}_1 \bar{\otimes} M_n)) \subset \text{CP}(\mathcal{M}_2 \bar{\otimes} M_n, \mathcal{N}_2 \bar{\otimes} M_n) \quad \forall n \in \mathbb{N}$$

Here, $I_n : \text{Hom}_{\mathbb{C}}(M_n, M_n) \rightarrow \text{Hom}_{\mathbb{C}}(M_n, M_n)$ is the identity map

In other words, S is completely positive if $S \otimes I_n$ preserves the linear ordering \preceq for all $n \in \mathbb{N}$

Properties of supermaps: Normality

A quantum supermap $S : \text{CB}(\mathcal{M}_1, \mathcal{N}_1) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{N}_2)$ must be continuous in a suitable sense

... But, if $\dim \mathcal{M}_i = \infty$ or $\dim \mathcal{N}_i = \infty$, many 'natural' topologies are available on the space $\text{CB}(\mathcal{M}, \mathcal{N})$ (e. g. pointwise uniform, strong, weak* or weak convergence. . .)

... So we avoid reference to a particular topology, and require normality with respect to a single suitable notion of increasing sequences in $\text{CB}(\mathcal{M}, \mathcal{N})$

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Increasing sequences in $\text{CB}(\mathcal{M}, \mathcal{N})$

Definition

A sequence $\{\mathcal{E}_n\}_{n \in \mathbb{N}}$ in $\text{CB}(\mathcal{M}, \mathcal{N})$ is

- *CP-increasing* if $0 \preceq \mathcal{E}_m \preceq \mathcal{E}_n$ for $m \leq n$
- *CP-bounded* if $\exists \mathcal{F} \in \text{CP}(\mathcal{M}, \mathcal{N})$ such that $\mathcal{E}_n \preceq \mathcal{F}$ for all n

Proposition

If the sequence $\{\mathcal{E}_n\}_{n \in \mathbb{N}}$ in $\text{CB}(\mathcal{M}, \mathcal{N})$ is CP-increasing and CP-bounded, then $\exists! \mathcal{E} \in \text{CP}(\mathcal{M}, \mathcal{N})$ such that

$$\text{wk}^*\text{-}\lim_{n \rightarrow \infty} \mathcal{E}_n(A) = \mathcal{E}(A) \quad \forall A \in \mathcal{M}$$

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$$\mathcal{E}_n \uparrow \mathcal{E}$$

Example (Kraus Theorem)

Suppose $\mathcal{M} \subset \mathcal{L}(\mathcal{H})$, $\mathcal{N} = \mathcal{L}(\mathcal{K})$. For all $\mathcal{E} \in \text{CP}(\mathcal{M}, \mathcal{L}(\mathcal{K})) \exists$ a sequence $\{E_k\}_{k \in \mathbb{N}}$ in $\mathcal{L}(\mathcal{K}, \mathcal{H})$ such that

$$\sum_{k=0}^n E_k^* \odot_{\mathcal{M}} E_k \uparrow \mathcal{E}$$

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Properties of supermaps: Normality

Definition

A linear map

$$S : \text{CB}(\mathcal{M}_1, \mathcal{N}_1) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{N}_2)$$

is *normal* if for all sequences $\{\mathcal{E}_n\}_{n \in \mathbb{N}}$ we have

$$\mathcal{E}_n \uparrow \mathcal{E} \quad \text{implies} \quad S(\mathcal{E}_n) \uparrow S(\mathcal{E})$$

Remark

A normal map

$$S : \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1)) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2))$$

is completely defined by its action on elementary tensors $E^* \odot_{\mathcal{M}_1} F$

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- (i) S is completely positive
- (ii) S is normal.

Definition

A quantum supermap S is *deterministic* if

$$S(\text{CP}_1(\mathcal{M}_1, \mathcal{N}_1)) \subset \text{CP}_1(\mathcal{M}_2, \mathcal{N}_2)$$

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The set of quantum supermaps is ordered: given two quantum supermaps S, T acting in the same CB spaces, we will write

$$S \ll T \quad \text{iff} \quad T - S \text{ is a quantum supermap}$$

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Example: Amplification

If \mathcal{V} is a Hilbert space, then

$$\begin{array}{ccc} \Pi_{\mathcal{V}} : \text{CB}(\mathcal{M}, \mathcal{N}) & \longrightarrow & \text{CB}(\mathcal{M} \bar{\otimes} \mathcal{L}(\mathcal{V}), \mathcal{N} \bar{\otimes} \mathcal{L}(\mathcal{V})) \\ \mathcal{E} & \longmapsto & \mathcal{E} \otimes \mathcal{I}_{\mathcal{V}} \end{array}$$

is a deterministic quantum supermap

Example: Concatenation

If $\mathcal{A} \in \text{CP}_0(\mathcal{N}_1, \mathcal{N}_2)$, $\mathcal{B} \in \text{CP}_0(\mathcal{M}_2, \mathcal{M}_1)$, then

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Theorem (Dilation theorem)

A linear map

$$S : \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1)) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2))$$

is a deterministic supermap iff there exist a triple $(\mathcal{V}, V, \mathcal{F})$, where

- \mathcal{V} is a Hilbert space
- $V : \mathcal{K}_2 \rightarrow \mathcal{K}_1 \otimes \mathcal{V}$ is an isometry
- $\mathcal{F} : \mathcal{M}_2 \rightarrow \mathcal{M}_1 \bar{\otimes} \mathcal{L}(\mathcal{V})$ is a quantum channel

such that

$$[S(\mathcal{E})](A) = V^* [(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})\mathcal{F}(A)] V$$

for all $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$ and $A \in \mathcal{M}_2$

Dilation of deterministic supermaps

- 1 The triple $(\mathcal{V}, V, \mathcal{F})$ can always be chosen in a way that

$$\mathcal{V} = \overline{\text{span}} \{ (u^* \otimes I_{\mathcal{V}}) V v \mid u \in \mathcal{K}_1, v \in \mathcal{K}_2 \}$$

(*minimal dilation*)

- 2 In this case, if $(\mathcal{V}', V', \mathcal{F}')$ is another dilation, then $\exists!$ isometry $W : \mathcal{V} \rightarrow \mathcal{V}'$ such that

$$V' = (I_{\mathcal{K}_1} \otimes W) V$$

and

$$\mathcal{F}(A) = (I_{\mathcal{M}_1} \otimes W^*) \mathcal{F}'(A) (I_{\mathcal{M}_1} \otimes W) \quad \forall A \in \mathcal{M}_2$$

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3 Setting

$$\mathcal{A} = V^* \odot_{\mathcal{L}(\mathcal{K}_1 \otimes \mathcal{V})} V$$

we have that S is the composition

$$S = C_{\mathcal{A}, \mathcal{F}} \circ \Pi_{\mathcal{V}}$$

4 In the Schrödinger (predual) picture

$$[S(\mathcal{E})]_*(\rho) = \mathcal{F}_*[(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})_*(V\rho V^*)]$$

for all $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$ and $\rho \in \mathcal{L}(\mathcal{K}_2)_* = \mathcal{T}(\mathcal{K}_2)$

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- 5 Taking $\mathcal{M}_1 = \mathcal{M}_2 = \mathbb{C}$, one gets Stinespring Theorem for normal CP maps
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Theorem (Radon-Nikodym theorem for supermaps)

Suppose $S, T : \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1)) \rightarrow \text{CB}(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2))$ are quantum supermaps, with

$$T \ll S.$$

Suppose S is deterministic, and let $(\mathcal{V}, V, \mathcal{F})$ be its minimal dilation. Then $\exists! \mathcal{G} \in \text{CP}_0(\mathcal{M}_2, \mathcal{M}_1 \bar{\otimes} \mathcal{L}(\mathcal{V}))$, with

$$\mathcal{G} \preceq \mathcal{F},$$

such that

$$[T(\mathcal{E})](A) = V^*[(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})\mathcal{G}(A)]V$$

for all $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$ and $A \in \mathcal{M}_2$

An example: Transforming POVMs into channels

- X : space of outcomes (discrete)
- $\mathcal{M}_1 = \ell^\infty(X)$: complex bounded functions (sequences) on X
- $\text{CP}_1(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$: $\mathcal{L}(\mathcal{K}_1)$ -valued POVMs on X

Let $(\mathcal{V}, V, \mathcal{F})$ be a dilation of a deterministic supermap

$$S : \text{CB}(\ell^\infty(X), \mathcal{L}(\mathcal{K}_1)) \longrightarrow \text{CB}(\mathcal{L}(\mathcal{H}_2), \mathcal{L}(\mathcal{K}_2))$$

For $x \in X$, set

$$\begin{aligned} \mathcal{F}_{x*} : \mathcal{T}(\mathcal{V}) &\longrightarrow \mathcal{T}(\mathcal{H}_2) \\ \sigma &\longmapsto \mathcal{F}_*(\delta_x \sigma) \end{aligned}$$

Then

$$[S(\mathcal{E})]_*(\rho) = \sum_{x \in X} \mathcal{F}_{x*}[(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})_*(V\rho V^*)_x]$$

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If $\mathcal{E} \in \text{CP}_1(\ell^\infty(X), \mathcal{L}(\mathcal{K}))$, then the map

$$\begin{aligned} P : X &\longrightarrow \mathcal{L}(\mathcal{K}) \\ x &\longmapsto \mathcal{E}(\delta_x) \end{aligned}$$

is a POVM, and

$$\mathcal{E}(f) = \sum_{x \in X} f_x P_x$$

$x \in X$

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Measure the POVM \mathcal{E} on \mathcal{K}_1 , thus obtaining the outcome $x \in X$

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Conditionally on the outcome x , apply the channel \mathcal{F}_{x*} on the ancilla \mathcal{V} , thus converting it into the output \mathcal{H}_2

- 1 Quantum operations
 - Definition
 - Dilation of quantum operations
- 2 Quantum supermaps
 - Definition
 - Dilation of quantum supermaps
 - Examples
- 3 Superinstruments
 - Definition
 - Dilation of superinstruments
 - Examples
- 4 Conclusions

Motivation

A quantum superinstrument R describes a measurement process of quantum channels

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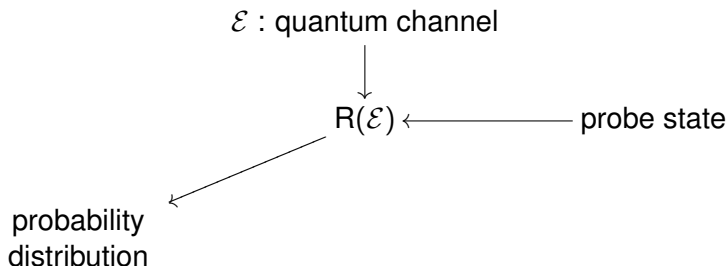
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\mathcal{E} : quantum channel



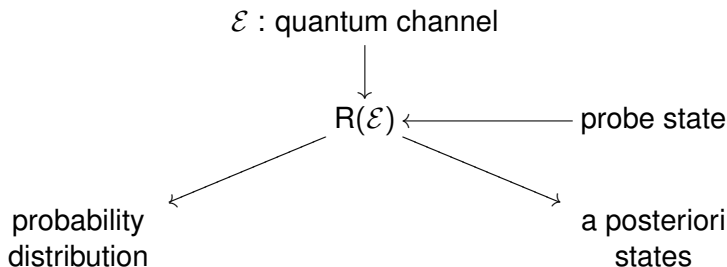
Motivation

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- $\text{CP}_1(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$: channels to be measured
- $\mathcal{L}(\mathcal{K}_2)_* = \mathcal{T}(\mathcal{K}_2)$: initial probe states
- \mathcal{M}_{2*} : final probe states
- (Ω, \mathcal{A}) : measurable set of outcomes
- $\mathcal{M}(\Omega; \text{CP}_0(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2)))$: quantum instruments on the probes

A quantum superinstrument is a map

$$R : \text{CP}_1(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1)) \longrightarrow \mathcal{M}(\Omega; \text{CP}_0(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2)))$$

To each channel $\mathcal{E} \in \text{CP}_1(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$ it associates the instrument

$$\mathcal{A} \ni B \longmapsto R_B(\mathcal{E}) \in \text{CP}_0(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2))$$

Motivation

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Recall that (in the Heisenberg picture) a classical quantum instrument is just a map

$$\mathcal{J} : \mathcal{A} \rightarrow \text{CP}_0(\mathcal{M}, \mathcal{L}(\mathcal{K}))$$

which is

- weak*-additive:

$$\mathcal{J}_B(A) = \text{wk}^*\text{-}\lim_{n \rightarrow \infty} \sum_{i=1}^n \mathcal{J}_{B_i}(A) \quad \forall A \in \mathcal{M} \quad \text{if } B_i \cap B_j = \emptyset \quad \forall i \neq j$$

- normalized:

$$\mathcal{J}_\Omega(I_{\mathcal{M}}) = I_{\mathcal{K}}$$

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Definition of superinstruments

Definition

Let (Ω, \mathcal{A}) be a measurable space. Suppose R is a map

$$\begin{array}{ccc} R : \mathcal{A} & \longrightarrow & \text{Hom}_{\mathbb{C}}(\text{CB}(\mathcal{M}_1, \mathcal{N}_1); \text{CB}(\mathcal{M}_2, \mathcal{N}_2)) \\ B & \longmapsto & R_B \end{array}$$

We say that R is a *quantum superinstrument* if

- (i) R_B is a quantum supermap for all $B \in \mathcal{A}$
- (ii) R_Ω is deterministic
- (iii) if $B = \bigcup_{i=1}^{\infty} B_i$ with $B_i \cap B_j = \emptyset$ for $i \neq j$, then

$$[R_B(\mathcal{E})](A) = \text{wk}^*\text{-}\lim_{n \rightarrow \infty} \sum_{i=1}^n [R_{B_i}(\mathcal{E})](A)$$

for all $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{N}_1)$ and $A \in \mathcal{M}_2$

Theorem (Dilation of quantum superinstruments)

Suppose that

$$R : \mathcal{A} \rightarrow \text{Hom}_{\mathbb{C}}(\text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1)); \text{CB}(\mathcal{M}_2, \mathcal{L}(\mathcal{K}_2)))$$

is a quantum superinstrument. Then there exist a Hilbert space \mathcal{V} , an isometry $V : \mathcal{K}_2 \rightarrow \mathcal{K}_1 \otimes \mathcal{V}$ and a quantum instrument

$$\mathcal{I} \in \mathcal{M}(\Omega; \text{CP}_0(\mathcal{M}_2, \mathcal{M}_1 \bar{\otimes} \mathcal{L}(\mathcal{V})))$$

such that

$$[R_B(\mathcal{E})](A) = V^*[(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})\mathcal{I}_B(A)]V \quad \forall A \in \mathcal{M}_2$$

for all $B \in \mathcal{A}$ and $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$

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$$\mathcal{I} \in \mathcal{M}(\Omega; \text{CP}_0(\mathcal{M}_2, \mathcal{M}_1 \bar{\otimes} \mathcal{L}(\mathcal{V})))$$

such that (in the Schrödinger picture)

$$[R_B(\mathcal{E})]_*(\rho) = \mathcal{I}_{B*}[(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})(V\rho V^*)] \quad \forall \rho \in \mathcal{T}(\mathcal{K}_2)$$

for all $B \in \mathcal{A}$ and $\mathcal{E} \in \text{CB}(\mathcal{M}_1, \mathcal{L}(\mathcal{K}_1))$

An example: Quantum testers

- $\text{CP}_1(\ell^\infty(X), \mathcal{L}(\mathcal{K}))$: input space of $\mathcal{L}(\mathcal{K})$ -valued POVMs on X
- \mathbb{C} : trivial output space

Fix a quantum superinstrument

$$R : \mathcal{A} \rightarrow \text{Hom}_{\mathbb{C}}(\text{CB}(\ell^\infty(X), \mathcal{L}(\mathcal{K})); \mathbb{C})$$

By dilation theorem

$$R_B(\mathcal{E}) = \langle v, (\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})(\mathcal{J}_B)v \rangle \quad \forall \mathcal{E} \in \text{CB}(\ell^\infty(X), \mathcal{L}(\mathcal{K})), B \in \mathcal{A}$$

For $x \in X$, define the $\mathcal{L}(\mathcal{V})$ -valued POVM on Ω

$$Q_x : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{V}) \quad Q_{x,B} = (\mathcal{J}_B)_x$$

Then

$$R_B(\mathcal{E}) = \sum_{x \in X} \text{tr} [Q_{x,B}(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})_*(\omega_v)_x]$$

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For $x \in X$, Q_x is a POVM on Ω

Ancillary Hilbert space \mathcal{V}

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$$R_B(\mathcal{E}) = \langle \mathbf{v}, (\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})(\mathcal{J}_B) \mathbf{v} \rangle \quad \forall \mathcal{E} \in \text{CB}(\ell^\infty(X), \mathcal{L}(\mathcal{K})), B \in \mathcal{A}$$

For $x \in X$ let $Q_{x,B} : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{V})$ be the POVM on \mathcal{V} defined by

Unit vector $\mathbf{v} \in \mathcal{K} \otimes \mathcal{V}$

$$Q_{x,B} = (\mathcal{J}_B)_x$$

Then

$$R_B(\mathcal{E}) = \sum_{x \in X} \text{tr} [Q_{x,B}(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})_*(\omega_{\mathbf{v}})_x]$$

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For $x \in$

Quantum instrument

$$\mathcal{J} : \mathcal{A} \longrightarrow \text{CP}(\mathbb{C}, \ell^\infty(X) \bar{\otimes} \mathcal{L}(\mathcal{V})) \simeq \ell^\infty(X; \mathcal{L}(\mathcal{V}))_+$$

Then

$$R_B(\mathcal{E}) = \sum_{x \in X} \text{tr} [Q_{x, B}(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})_*(\omega_v)_x]$$

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By dilation theorem (in the Schrödinger picture)

$$R_B(\mathcal{E}) = [\mathcal{I}_{B*}(\mathcal{E} \otimes \mathcal{I}_{\mathcal{V}})](\omega_{\mathcal{V}}) \quad \forall \mathcal{E} \in \text{CB}(\ell^\infty(X), \mathcal{L}(\mathcal{K})), B \in \mathcal{A}$$

Notation

$\omega_{\mathcal{V}}$: ortogonal projection on $\mathbb{C}v$ (rank-1 element in $\mathcal{T}(\mathcal{K} \otimes \mathcal{V})$)

Then

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Prepare a pure bipartite state $\omega_{\mathcal{V}}$ in $\mathcal{K} \otimes \mathcal{V}$

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Measure the POVM \mathcal{E} on \mathcal{K} , thus obtaining the outcome $x \in X$

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Conditionally on x , measure Q_x on \mathcal{V} , and obtain an outcome in B

- 1 Quantum operations
 - Definition
 - Dilation of quantum operations
- 2 Quantum supermaps
 - Definition
 - Dilation of quantum supermaps
 - Examples
- 3 Superinstruments
 - Definition
 - Dilation of superinstruments
 - Examples
- 4 Conclusions

Summary

- 1 We have given a general definition of supermaps
- 2 We have provided a dilation theorem for deterministic and probabilistic supermaps
- 3 We have characterized quantum superinstruments on the space of quantum channels
- 4 We have shown some applications

Open problems

- 1 How can the dilation theorems be extended to generic supermaps (not deterministic nor probabilistic)?
- 2 Is there a topology on $\text{CB}(\mathcal{M}, \mathcal{N})$ such that normality of supermaps is equivalent to continuity?

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- 3 Chiribella, G., D'Ariano, G. M., and Perinotti, P., A theoretical framework for quantum networks, *Phys. Rev. A* **80** (2009)

- ④ Chiribella, G., D'Ariano, G. M., and Perinotti, P., Quantum circuits architecture, Phys. Rev. Lett. **101**, 060401 (2008)
- ⑤ Chiribella, G., D'Ariano, G. M., and Perinotti, P., Transforming quantum operations: quantum supermaps, Europhys. Lett. **83**, 30004 (2008)
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