On the coexistence of conjugate observables on locally compact Abelian groups

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General results

- Weyl systems and conjugate observables
- Joint measurability for conjugate observables

Applications

- Continuous phase space
- Discrete phase space
- 3 Sketch of the main proof



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Weyl systems and conjugate observables

- G is a locally compact second countable Abelian group
- Ĝ is the dual of G
- μ and $\hat{\mu}$ are the Haar measures of G and \hat{G}
- $\langle \, \xi \, , \, x \,
 angle$ is the pairing of $\xi \in \hat{G}$ and $x \in G$

Definition (Weyl system)

Two unitary representations U of G and V of \hat{G} in the same Hilbert space \mathcal{H} form a *Weyl system* for the pair (G, \hat{G}) if

1) for all $x \in G$, $\xi \in \hat{G}$

$$U_X V_{\xi} = \overline{\langle \xi, x \rangle} V_{\xi} U_x;$$

2) there exists no nontrivial subspace $\mathcal{H}_0 \subset \mathcal{H}$ such that

 $U_x\mathcal{H}_0\subset\mathcal{H}_0$ and $V_{\xi}\mathcal{H}_0\subset\mathcal{H}_0$ $orall x\in G,\,\xi\in\hat{G}$.

Weyl systems and conjugate observables

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• for all $x \in G$, $\xi \in \hat{G}$

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 $U_x \mathcal{H}_0 \subset \mathcal{H}_0$ and $V_\xi \mathcal{H}_0 \subset \mathcal{H}_0$ $\forall x \in G, \, \xi \in \hat{G}$.

Definition (Conjugate observables)

Two observables $A : \mathcal{B}(G) \to \mathcal{L}(\mathcal{H})$ and $B : \mathcal{B}(\hat{G}) \to \mathcal{L}(\mathcal{H})$ are *conjugated* if

$$U_{X}A(X)U_{X}^{*} = A(X + x) \qquad V_{\xi}A(X)V_{\xi}^{*} = A(X)$$
$$U_{X}B(\Xi)U_{X}^{*} = B(\Xi) \qquad V_{\xi}B(\Xi)V_{\xi}^{*} = B(\Xi + \xi)$$
for all $x \in G, \xi \in \hat{G}, X \in \mathcal{B}(G), \Xi \in \mathcal{B}(\hat{G}).$

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By Stone-Naimark-Ambrose-Godement Theorem,

$$U_{\mathbf{x}} = \int_{\hat{\mathbf{G}}} \overline{\langle \xi, \mathbf{x} \rangle} \, \mathrm{d}\hat{\mathsf{B}}(\xi) \qquad V_{\xi} = \int_{\mathbf{G}} \langle \xi, \mathbf{x} \rangle \, \mathrm{d}\hat{\mathsf{A}}(\xi) \, .$$

where \hat{A} and \hat{B} are conjugated spectral measures.

Every Weyl system is unitarily equivalent to the canonical one

$$\mathcal{H} = L^2(G, \mu)$$
$$U_X f = f(\cdot - x) \qquad V_{\xi} f = \langle \xi, \cdot \rangle f$$
$$\hat{A}(X) f = \mathbb{1}_X f \qquad \hat{B}(\Xi) f = (\mathcal{F}^{-1} \mathbb{1}_{\Xi}) * f$$

where $\mathcal{F} : L^2(G, \mu) \to L^2(\hat{G}, \hat{\mu})$ is Fourier transform and * is convolution.

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where $\mathcal{F} : L^2(G, \mu) \to L^2(\hat{G}, \hat{\mu})$ is Fourier transform and * is convolution.

$$\mathcal{H} = L^{2}(\mathbb{R}, \, \mathrm{d}x)$$
$$U_{x} = \mathrm{e}^{-ixP} \qquad V_{\xi} = \mathrm{e}^{i\xi Q}$$
$$P = \int_{\mathbb{R}} z \, \mathrm{d}\hat{\mathsf{A}}(z) \qquad Q = \int_{\mathbb{R}} z \, \mathrm{d}\hat{\mathsf{B}}(z)$$

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$$\mathcal{H} = \mathbb{C}^d$$

 $U_x e_k = e_{k+x}$ $V_\xi e_k = e^{2\pi i \frac{\xi k}{d}} e_k$
 $\hat{A}(k) = |e_k\rangle \langle e_k|$ $\hat{B}(k) = \mathcal{F}^* \hat{A}(k) \mathcal{F}$.

In particular, the observavles \hat{A} and \hat{B} are *mutually unbiased* tr $\left[\hat{A}(i)\hat{B}(j)\right] = 1/d \quad \forall i, j$.

Theorem (Carmeli, Heinosaari, T.)

If (A, B) are conjugate observables, there exist unique probability measures Λ on G and Γ on \hat{G} such that

$$\begin{aligned} \mathsf{A}(X) &=: \mathsf{A}_{\Lambda}(X) = \int_{G} \Lambda(X - x) \, d\hat{\mathsf{A}}(x) \qquad \forall X \in \mathcal{B}(G) \\ \mathsf{B}(\Xi) &=: \mathsf{B}_{\Gamma}(\Xi) = \int_{\hat{G}} \Gamma(\Xi - \xi) \, d\hat{\mathsf{B}}(\xi) \qquad \forall \Xi \in \mathcal{B}(\hat{G}) \, . \end{aligned}$$

The observables $(A_{\Lambda}, B_{\Gamma})$ are *noisy versions* of (\hat{A}, \hat{B}) :

 $\begin{aligned} p_{\rho}^{\mathsf{A}_{\Lambda}}(X) &:= \operatorname{tr}\left[\rho\mathsf{A}_{\Lambda}(X)\right] = (\Lambda * p_{\rho}^{\hat{\mathsf{A}}})(X) \\ p_{\rho}^{\mathsf{B}_{\Gamma}}(\Xi) &:= \operatorname{tr}\left[\rho\mathsf{B}_{\Gamma}(\Xi)\right] = (\Gamma * p_{\rho}^{\hat{\mathsf{B}}})(\Xi) \end{aligned} \quad \forall \rho \in \mathcal{S}(\mathcal{H}) . \end{aligned}$

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Definition (Covariant phase-space observables)

An observable G : $\mathcal{B}(G \times \hat{G}) \rightarrow \mathcal{L}(\mathcal{H})$ is a *covariant phase-space observable* if $U_x V_{\xi} G(Z) V_{\xi}^* U_x^* = G(Z + (x, \xi))$ for all $Z \in \mathcal{B}(G \times \hat{G}), x \in G, \xi \in \hat{G}$.

Theorem (Holevo; Cassinelli, De Vito, T.)

If G is a covariant phase-space observable, then there exists a unique state $\tau \in S(\mathcal{H})$ such that

$$\mathsf{G}(Z) = \int_{Z} U_{\mathsf{X}} V_{\xi} \tau \, V_{\xi}^* \, U_{\mathsf{X}}^* \, d\mu(\mathsf{X}) \, d\hat{\mu}(\xi) \qquad \forall Z \in \mathcal{B}(G \times \hat{G}) \,.$$

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Taking the margins of a covariant phase-space observable

$$A(X) = G(X \times \hat{G})$$
 $B(\Xi) = G(G \times \Xi)$

we obtain conjugate observables (A, B).

Definition (Joint measurability)

Conjugate observables (A, B) are *jointly measurabile* if they are the margins of a single observable G on $G \times \hat{G}$.

Remark

The joint observable G is in general not unique.

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Main theorem

Taking the margins of a covariant phase-space observable

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Theorem (Carmeli, Heinosaari, T.)

Conjugate observables (A, B) are jointly measurabile if and only if they are the margins of a covariant phase-space observable G on $G \times \hat{G}$.

Remark

The generating state τ of the covariant joint observable $G=G_{\tau}$ is in general not unique.

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Corollaries

Corollary

For conjugate observables $(A_\Lambda,B_\Gamma),$ these are equivalent facts:

- $(A_{\Lambda}, B_{\Gamma})$ are jointly measurable;
- there is a state $au \in \mathcal{S}(\mathcal{H})$ such that

$$\Lambda(X) = \operatorname{tr}\left[\tau \hat{\mathsf{A}}(-X)\right] \qquad \Gamma(\Xi) = \operatorname{tr}\left[\tau \hat{\mathsf{B}}(-\Xi)\right];$$

- there is a vector $\psi \in \mathcal{H} \otimes \mathcal{H}$ such that

$$\Lambda(X) = \left\langle \psi, \, (\hat{\mathsf{A}}(-X) \otimes \mathbb{1})\psi \right\rangle \qquad \Gamma(\Xi) = \left\langle \psi, \, (\hat{\mathsf{B}}(-X) \otimes \mathbb{1})\psi \right\rangle.$$

Corollary

Conjugate observables $(A_{\Lambda}, B_{\Gamma})$ are jointly measurable only if Λ and Γ have densities w.r.t. the Haar measures μ and $\hat{\mu}$.

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If p is a probability measure on \mathbb{R} ,

$$\operatorname{Var}(p) = \int_{\mathbb{R}} \left(y - \int_{\mathbb{R}} x \, \mathrm{d}p(x) \right)^2 \, \mathrm{d}p(y) \, .$$

Proposition

If the conjugate observables (A, B) are jointly measurable, then for all states $\rho \in S(\mathcal{H})$ $\operatorname{Var}(p^{A}_{\circ})\operatorname{Var}(p^{B}_{\circ}) > 1$.

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Case $G = \hat{G} = \mathbb{Z}_d$

If $\lambda,\gamma\in[0,1],$ we define conjugate observables $(\mathsf{A}_{\lambda},\mathsf{B}_{\gamma})$

$$\begin{split} \mathsf{A}_{\lambda}(k) &= \lambda \hat{\mathsf{A}}(k) + (1-\lambda) \frac{1}{d}\mathbb{1} \\ \mathsf{B}_{\gamma}(k) &= \gamma \hat{\mathsf{B}}(k) + (1-\gamma) \frac{1}{d}\mathbb{1} \,. \end{split}$$

Proposition

For all $\lambda \in [0, 1]$, let

$$\gamma_{\max}(\lambda) = rac{(d-2)(1-\lambda)+2\sqrt{(1-d)\lambda^2+(d-2)\lambda+1}}{d}$$

Then, the conjugate observables $(A_{\lambda}, B_{\gamma})$ are jointly measurable if and only if $0 \leq \gamma \leq \gamma_{max}(\lambda)$.

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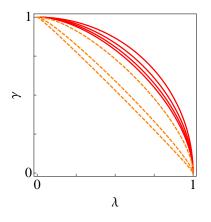


Figure: The boundary curve $\lambda \mapsto \gamma_{max}(\lambda)$ for d = 2, 3, 4, 5 (red solid curves) and for d = 10, 100, 1000 (orange dashed curves).

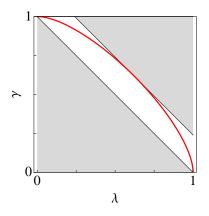


Figure: In this picture d = 10. The grey regions represent necessary and sufficient linear conditions.

Corollary

Conjugate observables $(A_{\lambda}, B_{\gamma})$

- are jointly measurable if $\gamma + \lambda \leq$ 1;
- are not jointly measurable if $\gamma + \lambda \ge 1 + \frac{\sqrt{d-1}}{d-1}$.

Proposition

If $\gamma = \gamma_{max}(\lambda)$, then the conjugate observables $(A_{\lambda}, B_{\gamma})$ have a unique joint observable. This unique joint observable is the covariant phase-space observable G_{τ} with

$$au = |\chi_{\lambda}\rangle \langle \chi_{\lambda}| , \qquad \chi_{\lambda} = \alpha_{\lambda} e_{0} + \beta_{\lambda} \mathcal{F}^{*} e_{0} ,$$

where

$$\alpha_{\lambda} = \frac{1}{\sqrt{d}} \left[\sqrt{(d-1)\lambda + 1} - \sqrt{1-\lambda} \right] \qquad \beta_{\lambda} = \sqrt{1-\lambda} \,.$$

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$$\begin{split} \mathsf{A}_{\lambda;p}(j) &= \lambda \mathsf{A}(j) + (1-\lambda)p(j)\mathbb{1} \\ \mathsf{B}_{\gamma;r}(k) &= \gamma \mathsf{B}(k) + (1-\gamma)r(k)\mathbb{1} \,. \end{split}$$

The observables $A_{\lambda;p}$ and $B_{\gamma;r}$ are *not* covariant. However, the following still holds.

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Sequential implementation

Define the A_{λ} -compatible instrument

$$\begin{aligned} \mathcal{I} : \mathbb{Z}_{d} \times \mathcal{S}(\mathcal{H}) \to \mathcal{S}(\mathcal{H}) \\ \mathcal{I}(k, \rho) = \sqrt{\mathsf{A}_{\lambda}(k)} \rho \sqrt{\mathsf{A}_{\lambda}(k)} \end{aligned}$$

Then, \mathcal{I} is a generalized position instrument

 $U_i V_j \mathcal{I}(k, V_j^* U_i^* \rho U_i V_j) V_j^* U_i^* = \mathcal{I}(i+k, \rho)$

and a sequential measurement of ${\mathcal I}$ followed by B gives the covariant phase-space observable

$$\mathsf{G}(i,j) = \mathcal{I}_i^*(\mathsf{B}(j)) = \sqrt{\mathsf{A}_{\lambda}(i)}\mathsf{B}_{\gamma}(j)\sqrt{\mathsf{A}_{\lambda}(i)}\,.$$

It is easy to check that $G \equiv G_{\tau}$, with $\tau = |\chi_{\lambda}\rangle \langle \chi_{\lambda}|$. Therefore, G is a joint observable of $(A_{\lambda}, B_{\gamma})$ with $\gamma \equiv \gamma_{max}(\lambda)$.

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Proposition

Suppose $(A_{\lambda}, B_{\gamma})$ are conjugate observables.

- If λ ∉ {0,1} and 0 < γ < γ_{max}(λ), then they have an informationally complete joint observable.
- If $\lambda \notin \{0, 1\}$ and $\gamma = \gamma_{\max}(\lambda)$, then their unique joint observable is informationally complete if and only if d is odd.
- If $\lambda = 0$ or $\gamma = 0$, then they have no informationally complete joint observable.

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General results

- Weyl systems and conjugate observables
- Joint measurability for conjugate observables

Applications

- Continuous phase space
- Discrete phase space

3 Sketch of the main proof

4 References

Case $|G| = |\hat{G}| < \infty$.

If conjugate observables (A, B) are jointly measurable, with joint observable G on $G \times \hat{G}$, we define

$$\tilde{\mathsf{G}}(x,\xi) = \frac{1}{|G|^2} \sum_{(y,\eta)\in G\times\hat{G}} \underbrace{U_y^* V_\eta^* \mathsf{G}(x+y,\xi+\eta) V_\eta U_y}_{\mathsf{G}_{(y,\eta)}(x,\xi)} \,.$$

Then, G is a covariant phase-space observable with margins

$$\tilde{\mathsf{G}}(\{x\} \times \hat{G}) = \frac{1}{|G|^2} \sum_{(y,\eta) \in G \times \hat{G}} U_y^* V_\eta^* \sum_{\xi \in \hat{G}} \mathsf{G}(x+y,\xi+\eta) V_\eta U_y$$
$$= \frac{1}{|G|^2} \sum_{(y,\eta) \in G \times \hat{G}} U_y^* V_\eta^* \mathsf{A}(x+y) V_\eta U_y = \mathsf{A}(x) \,.$$

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$$\begin{split} \tilde{\mathsf{G}}(\boldsymbol{G} \times \{\xi\}) &= \frac{1}{|\boldsymbol{G}|^2} \sum_{(\boldsymbol{y}, \eta) \in \boldsymbol{G} \times \hat{\boldsymbol{G}}} \boldsymbol{U}_{\boldsymbol{y}}^* \boldsymbol{V}_{\boldsymbol{\eta}}^* \sum_{\boldsymbol{x} \in \boldsymbol{G}} \mathsf{G}(\boldsymbol{x} + \boldsymbol{y}, \xi + \boldsymbol{\eta}) \boldsymbol{V}_{\boldsymbol{\eta}} \boldsymbol{U}_{\boldsymbol{y}} \\ &= \frac{1}{|\boldsymbol{G}|^2} \sum_{(\boldsymbol{y}, \eta) \in \boldsymbol{G} \times \hat{\boldsymbol{G}}} \boldsymbol{U}_{\boldsymbol{y}}^* \boldsymbol{V}_{\boldsymbol{\eta}}^* \mathsf{B}(\xi + \boldsymbol{\eta}) \boldsymbol{V}_{\boldsymbol{\eta}} \boldsymbol{U}_{\boldsymbol{y}} = \mathsf{B}(\xi) \end{split}$$

Case $|G| = |\hat{G}| = \infty$

Definition (Invariant mean)

Let Ω be a locally compact separable metric space. An *operator valued mean* on Ω is a linear map

 $\mathsf{M}:\textit{BC}(\Omega)\to\mathcal{L}(\mathcal{K})$

such that

(i)
$$M(f) \ge 0$$
 if $f \ge 0$;
(ii) $M(1) = 1$.

An observable M on Ω defines a operator valued mean on Ω

$$\mathsf{M}(f) := \int_{\Omega} f(\omega) \, \mathrm{d} \mathsf{M}(\omega) \qquad \forall f \in BC(\Omega) \,.$$

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Remark

If $\mathcal{K} = \mathbb{C}$, then M is a *mean*, that is, a normalized positive linear functional on $BC(\Omega)$.

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$\mathsf{M}(\infty) = \mathbb{1} - \mathrm{LUB}\{\mathsf{M}(f) \mid f \in C_{\mathcal{C}}(\Omega), \ \mathbf{0} \leq f \leq 1\}.$

There exists a unique operator valued measure M_0 on Ω such that $M(f) = M_0(f) \qquad \forall f \in C_c(\Omega)$.

For such M₀,

$$\mathsf{M}_0(\Omega) = \mathbb{1} - \mathsf{M}(\infty) \,.$$

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If $M(\infty) = 0$, then

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Alessandro Toigo (PoliMi, INFN) Coexistence of conjugate observables 25/31 Munich, September 10th, 2013 25 / 31

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If M is an operator valued mean on Ω , then the *i*-th margin of M is the operator valued mean M_i on Ω_i given by

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A mean *m* on $G \times \hat{G}$ is *invariant* if

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Theorem (E. Hewitt, K. A. Ross, Th. IV.17.5)

There exists an invariant mean on $G \times \hat{G}$.

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Let M be a joint observable of A_{Λ} and B_{Γ} . For all $\varphi, \psi \in \mathcal{H}$ and $(x, \xi) \in G \times \hat{G}$, define $\Theta[f; \varphi, \psi](x, \xi) = \langle \mathsf{M}(f_{(x,\xi)}) V_{\xi}^* U_x^* \varphi, V_{\xi}^* U_x^* \psi \rangle$. Then, $\Theta[f; \varphi, \psi] \in BC(G \times \hat{G})$, and $\Theta[f_{(x,\xi)}; \varphi, \psi] = \Theta[f; U_x V_{\xi} \varphi, U_x V_{\xi} \psi]_{(x,\xi)}$ $\Theta[\tilde{f}_1; \varphi, \psi](x, \xi) = \langle \mathsf{A}_{\Lambda}(f)\varphi, \psi \rangle \quad \forall f \in BC(G)$ $\Theta[\tilde{q}_2; \varphi, \psi](x, \xi) = \langle \mathsf{B}_{\Gamma}(q)\varphi, \psi \rangle \quad \forall a \in BC(\hat{G})$.

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$$egin{array}{lll} \mathsf{M}^{\mathrm{av}}: \textit{BC}(\textit{G} imes \hat{\textit{G}})
ightarrow \mathcal{L}(\mathcal{H}) \ \langle \, \mathsf{M}^{\mathrm{av}}(\textit{f}) arphi \,, \, \psi \,
angle = \textit{m}(\Theta[\textit{f}; arphi, \psi]) \,. \end{array}$$

Then, M^{av} is a operator valued mean on $G \times \hat{G}$ satisfying $M^{av}(f_{(x, c)}) = V^*_{c} U^*_{x} M^{av}(f) U_x V_{c}.$

$$\mathsf{M}^{\mathrm{av}}(f_{(x,\xi)}) = V_{\xi}^* U_x^* \mathsf{M}^{\mathrm{av}}(f) U_x V_{\xi}$$

Moreover,

$$\begin{array}{ccc} M_1^{\mathrm{av}} = A_\Lambda \\ M_2^{\mathrm{av}} = B_\Gamma \end{array} & \Longrightarrow & \begin{array}{c} M_1^{\mathrm{av}}(\infty) = 0 \\ M_2^{\mathrm{av}}(\infty) = 0 \end{array} \implies & M^{\mathrm{av}}(\infty) = 0 \,. \end{array}$$

Therefore, $M^{av} = M^{av}_0$, that is, M^{av} is an operator valued measure on $G \times \hat{G}$.

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Fix an invariant mean *m* on $G \times \hat{G}$, and define the map

$$egin{aligned} \mathsf{M}^{\mathrm{av}} &: \textit{BC}(\textit{G} imes \hat{\textit{G}})
ightarrow \mathcal{L}(\mathcal{H}) \ \langle \, \mathsf{M}^{\mathrm{av}}(\textit{f}) arphi \,, \, \psi \,
angle &= \textit{m}(\Theta[\textit{f}; arphi, \psi]) \,. \end{aligned}$$

Then, M^{av} is a operator valued mean on $G \times \hat{G}$ satisfying

$$\mathsf{M}^{\mathrm{av}}(f_{(x,\xi)}) = V_{\xi}^* U_x^* \mathsf{M}^{\mathrm{av}}(f) U_x V_{\xi} \,.$$

Moreover,

$$\begin{array}{ccc} M_1^{\rm av}=A_\Lambda & & & M_1^{\rm av}(\infty)=0 \\ M_2^{\rm av}=B_\Gamma & & & M_2^{\rm av}(\infty)=0 \end{array} \implies & M^{\rm av}(\infty)=0 \ . \end{array}$$

Therefore, $M^{av} = M^{av}_0$, that is, M^{av} is an operator valued measure on $G \times \hat{G}$.

General results

- Weyl systems and conjugate observables
- Joint measurability for conjugate observables

Applications

- Continuous phase space
- Discrete phase space
- 3 Sketch of the main proof



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